

# Update

**Jose A. Vazquez**

BNL Meeting  
November,

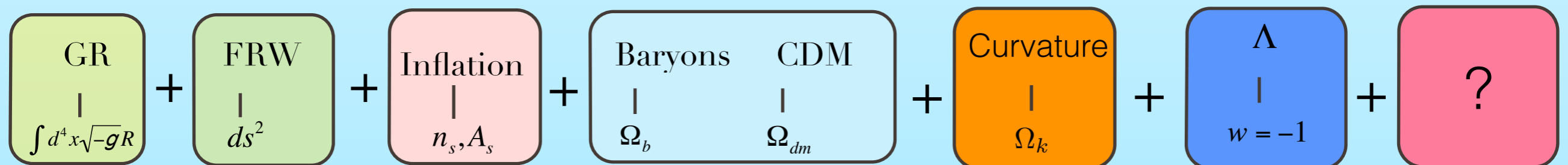
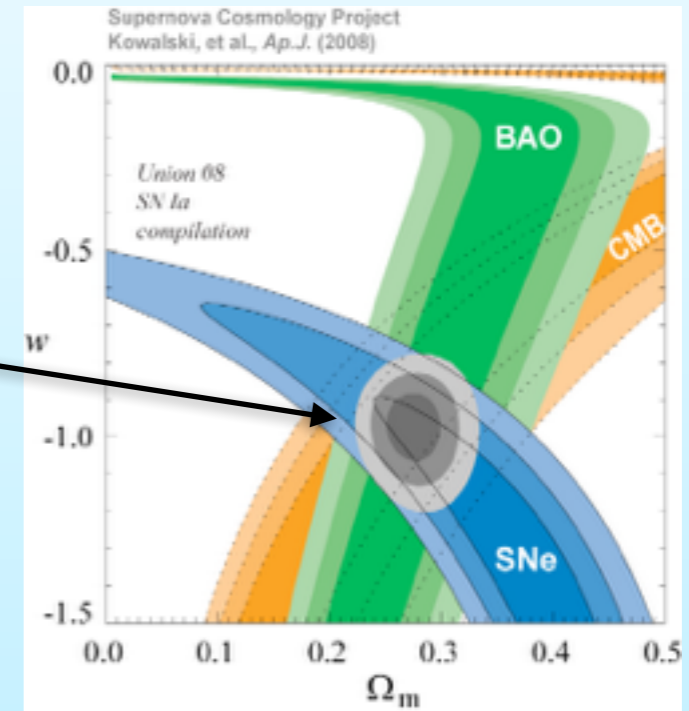
# Early Dark Energy

DR12

Cosmological parameter constraints from galaxy-galaxy  
lensing and galaxy clustering

BNL Meeting  
November,

# The standard cosmological model



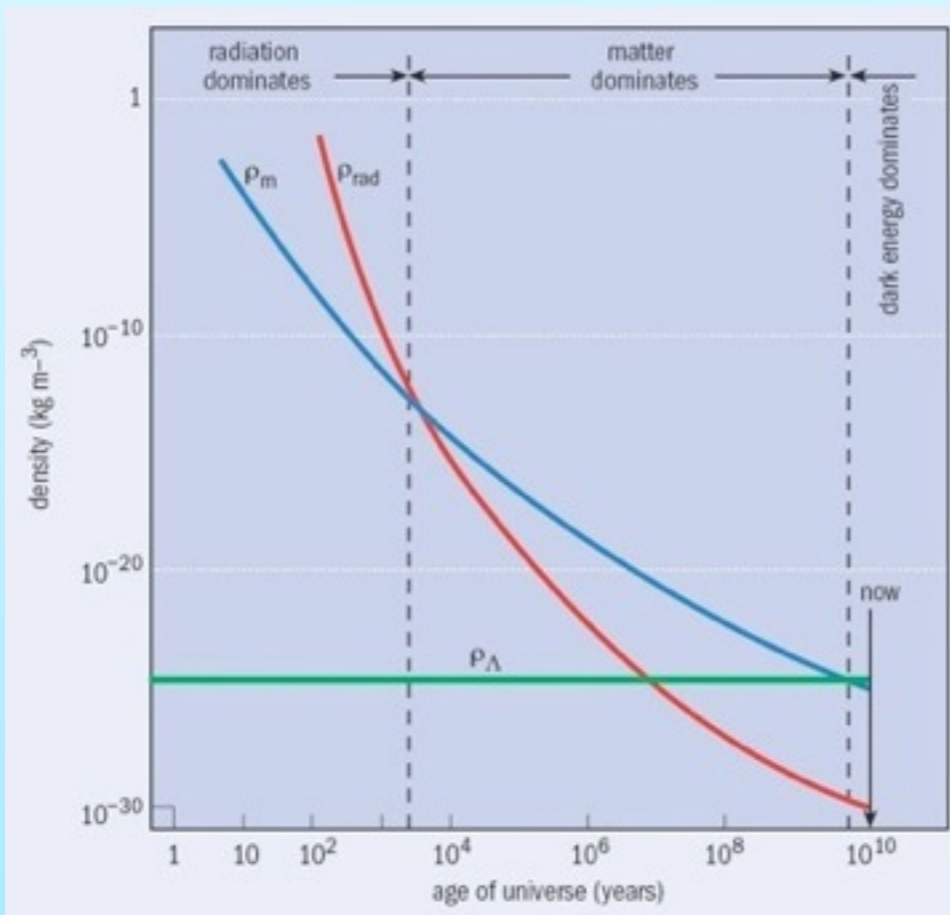
# Cosmological constant

cosmology vs particle physics:

$$\langle \rho \rangle \approx 2^{-10} \pi^{-4} G^{-2} = 2 \times 10^{71} \text{ GeV}^4 .$$

$$|\rho_V| \lesssim 10^{-29} \text{ g/cm}^3 \approx 10^{-47} \text{ GeV}^4 .$$

clearly a discrepancy!!



fine tuning problem

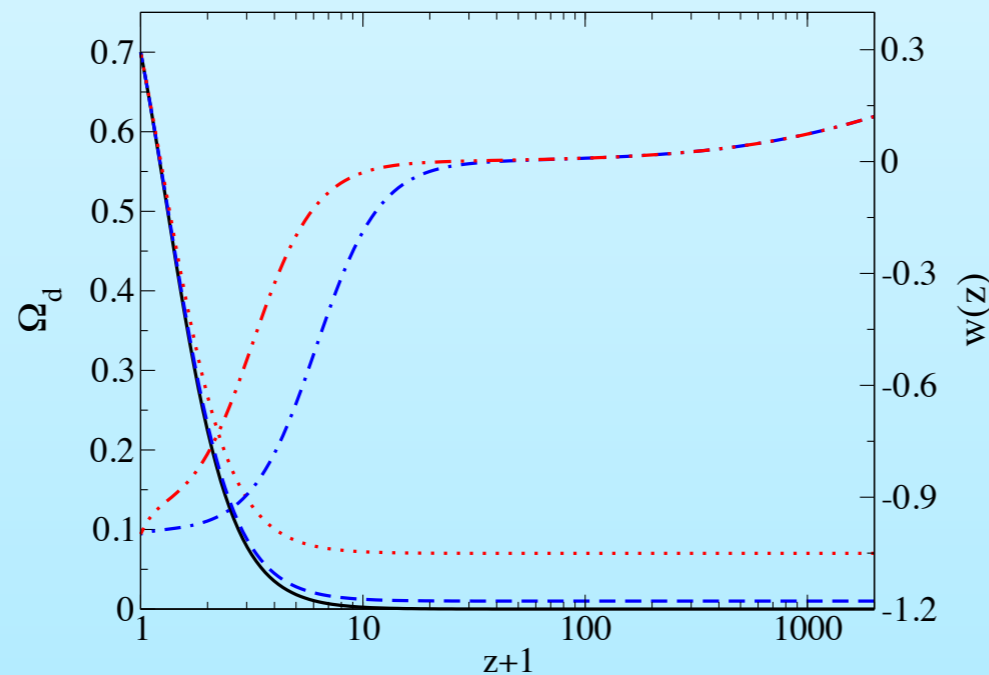
coincidence problem

# Early Dark Energy

In typical dark energy models, **DE is dynamically negligible** at high redshifts.

Some scalar field potentials **track the energy density** of the dominant species **during the radiation and matter** dominated eras.

$$\Omega_{\text{de}}(a) = \frac{\Omega_{\text{de}} - \Omega_{\text{de}}^e (1 - a^{-3w_0})}{\Omega_{\text{de}} + \Omega_m a^{3w_0}} + \Omega_{\text{de}}^e (1 - a^{-3w_0})$$



$$\Omega_{\text{de}}^e$$

- To describe the BAO peak;

Across the line of sight

$$D_C(z) = \frac{c}{H_0} \int_0^z dz' \frac{H_0}{H(z')}$$

Along the line of sight

$$D_H(z) = c/H(z)$$

The length of the ruler:

$$r_d = \int_{z_d}^{\infty} \frac{c_s(z)}{H(z)} dz$$

$$D_M(z)/r_d = \alpha_{\perp} D_{M,\text{fid}}(z)/r_{d,\text{fid}}$$

$$D_H(z)/r_d = \alpha_{\parallel} D_{H,\text{fid}}(z)/r_{d,\text{fid}}$$



**The BAO signal**

# Some approximations

$$\rho_m(z) + \rho_{de}(z) = 3H^2(z)/8\pi G$$

$$H(z) = (100 \text{ km s}^{-1} \text{ Mpc}^{-1}) \times [(\Omega_m h^2) (1+z)^3]^{1/2} \times [1 + \rho_{de}(z)/\rho_m(z)]^{1/2}.$$



$$\Omega_{de}^e / \Omega_m(z) \approx \Omega_{de}^e / (1 - \Omega_{de}^e)$$

at high redshift

•  $H(z)$

is higher in the early dark energy model by a factor  $(1 - \Omega_{de}^e)^{-1/2}$ , and  $D_H(z)$  is smaller by the same factor.

the boosted energy density in this era reduces

•  $r_d$  sound horizon by a factor  $(1 - \Omega_{de}^e)^{1/2}$

•  $D_H(z)$  is smaller by the same factor

$\alpha_{||}$

at low redshift

$$D_M(z) = \frac{c}{H_0} \int_0^z \frac{H_0}{H(z')} dz'$$

depends mainly on  $H_0$

Therefore, to keep the value of  $D_M(1090)/r_d$  fixed to the CMB constraint, one must increase  $H_0$  by approximately  $(1 - \Omega_{de}^e)^{-1/2}$

$\alpha_{\perp}$

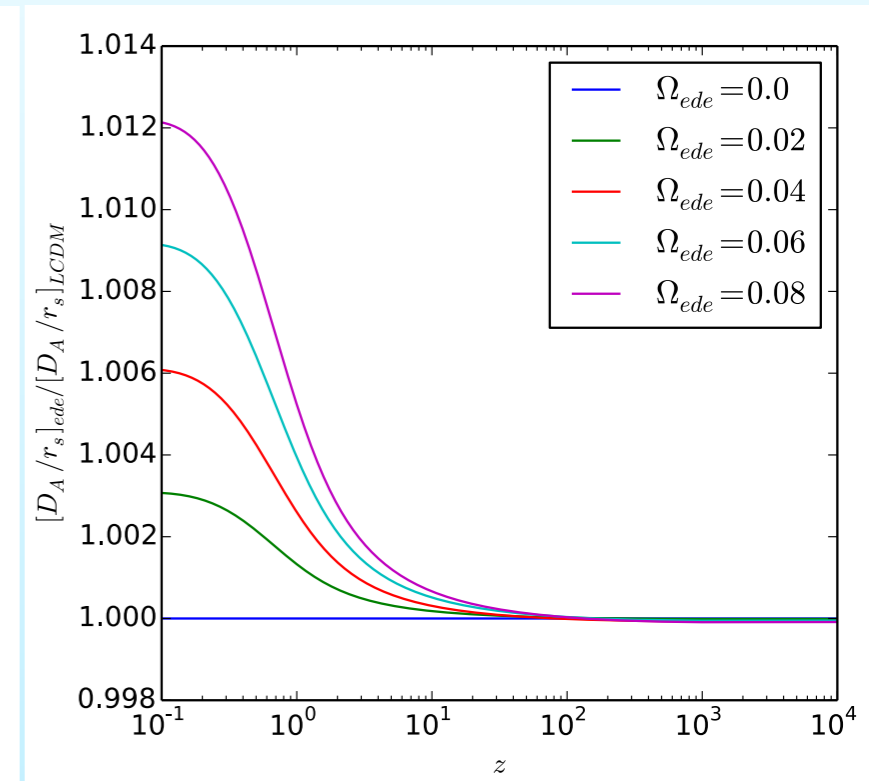
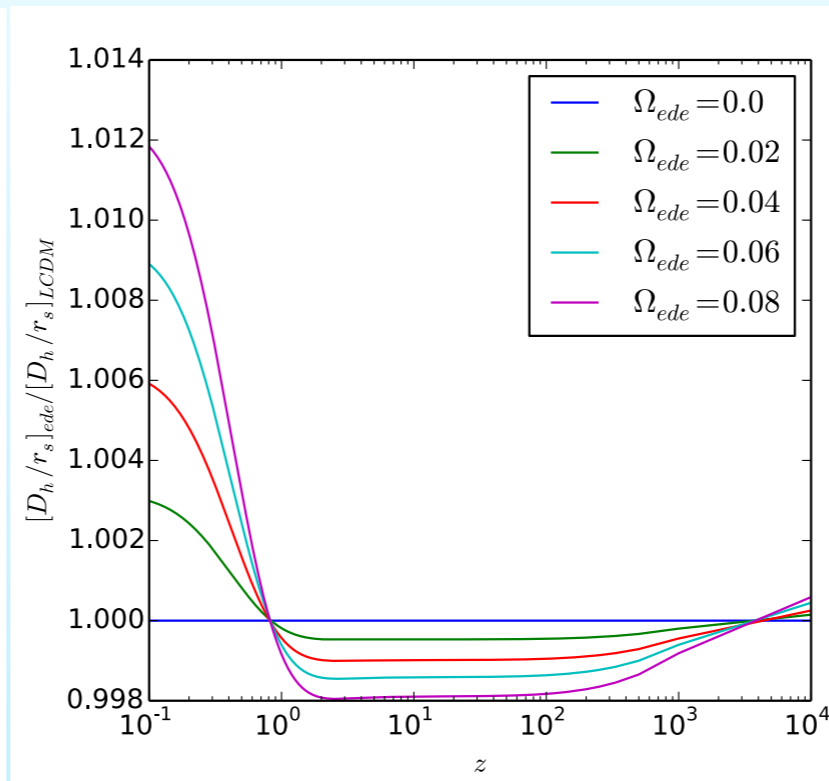
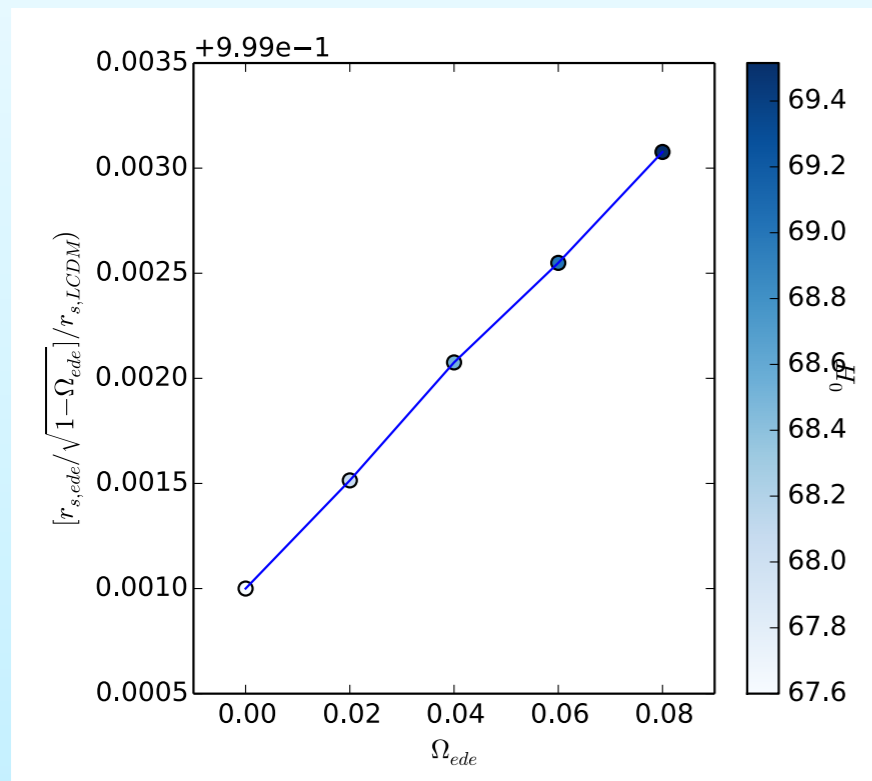
# Early Dark Energy

$r_d$

sound horizon by a factor  $(1 - \Omega_{de}^e)^{1/2}$

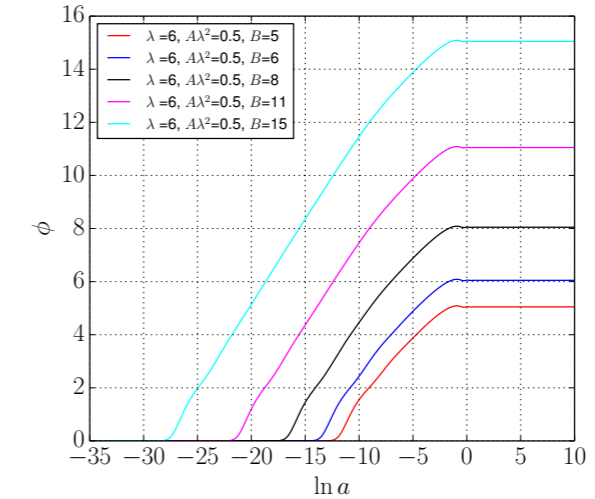
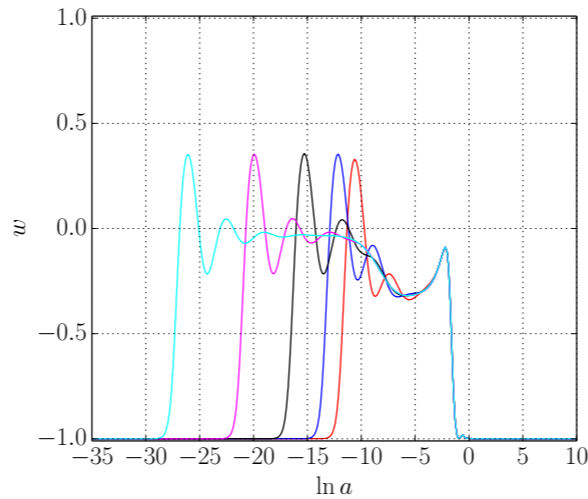
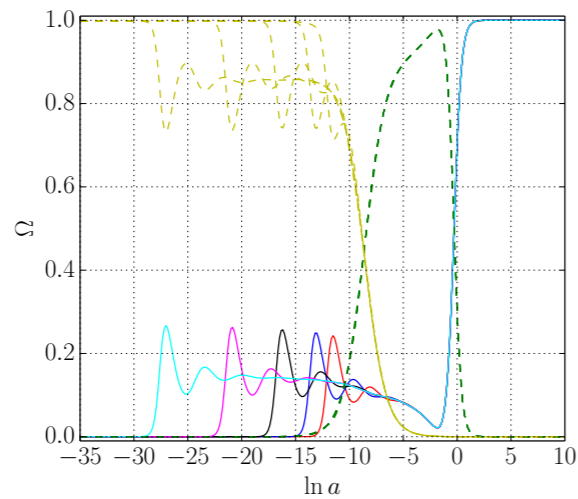
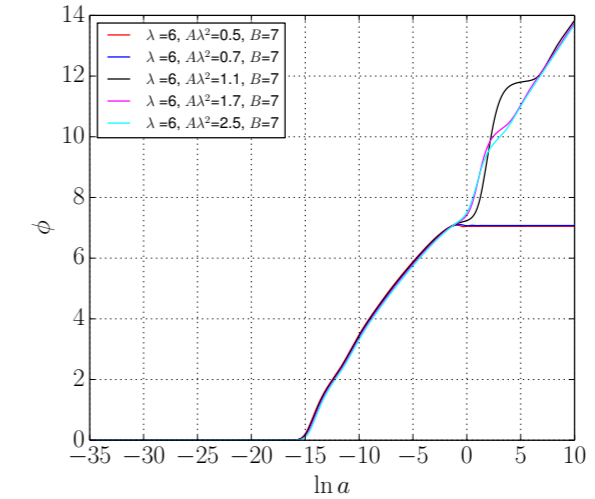
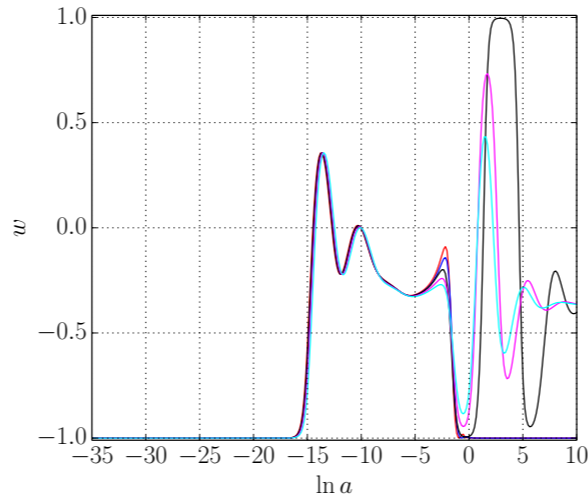
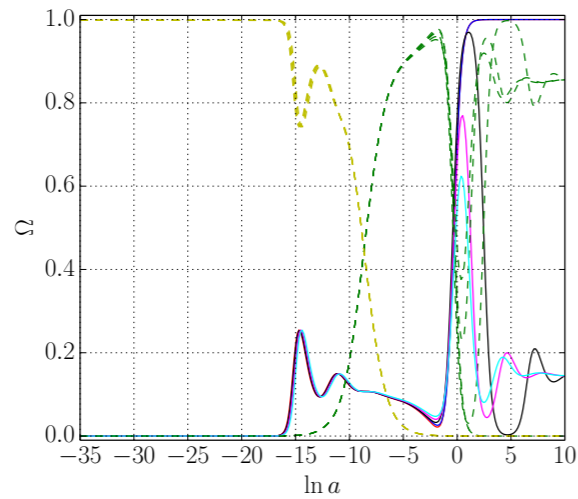
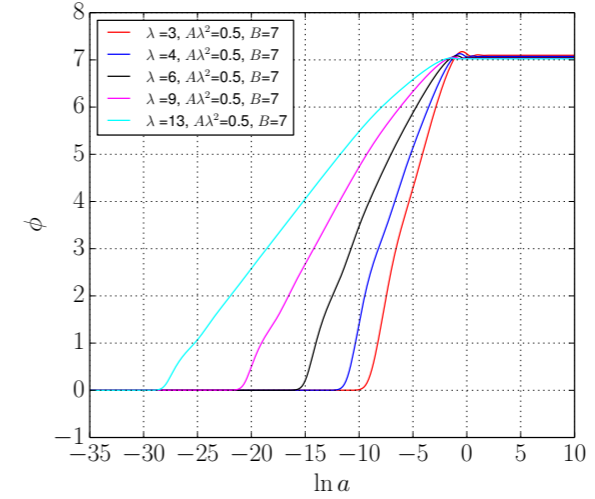
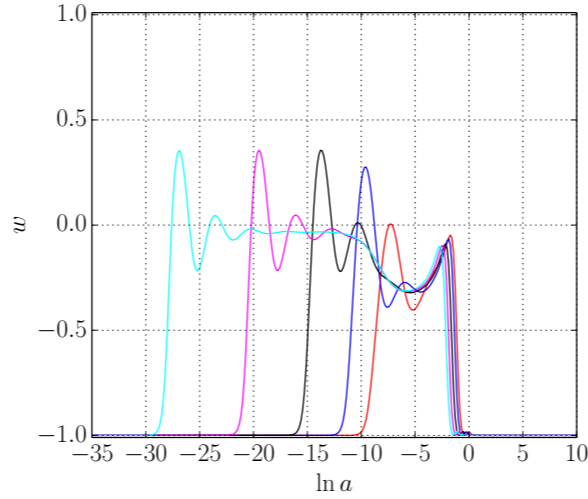
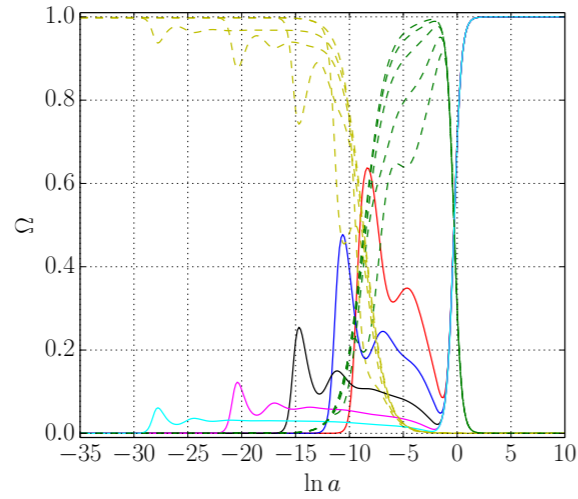
$\alpha_{||}$

$\alpha_{\perp}$

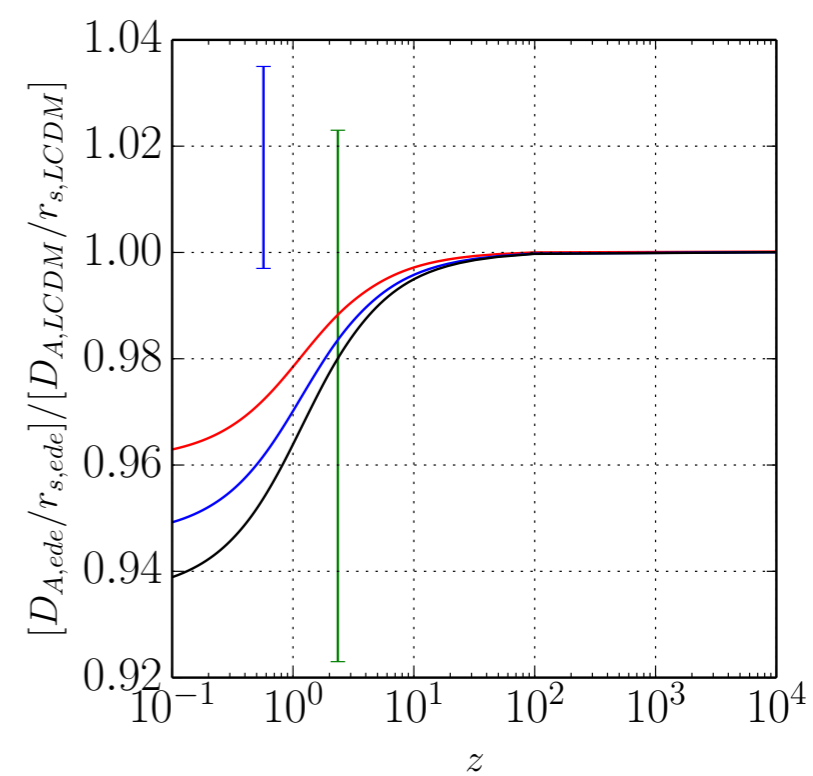
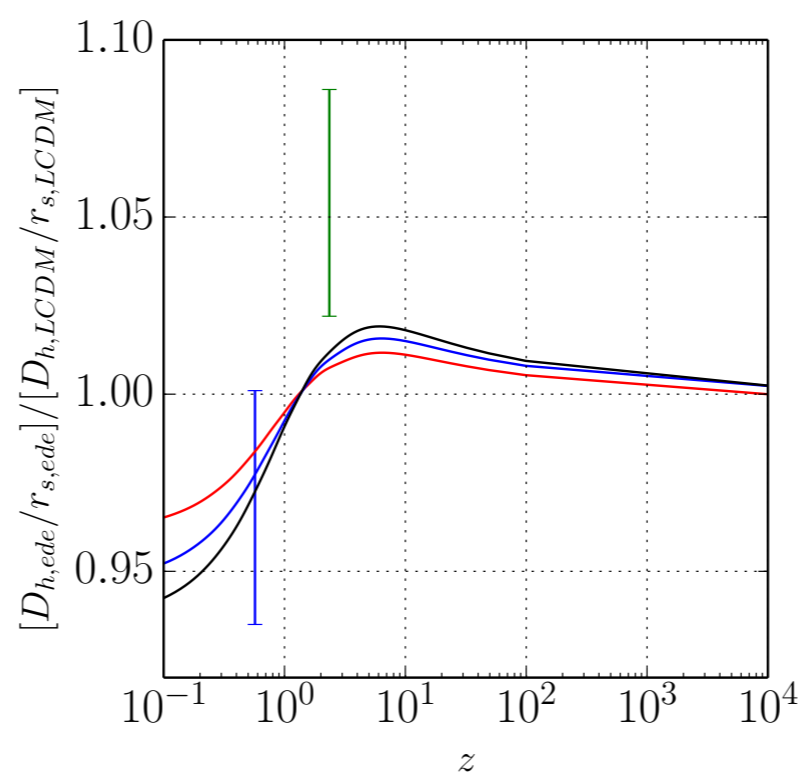
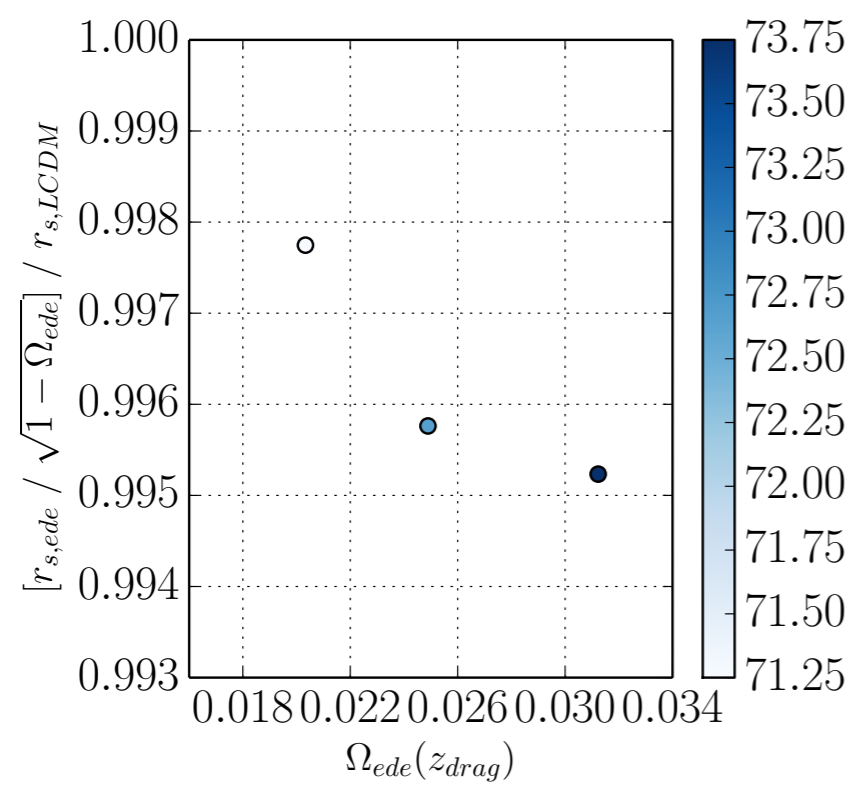


# Scalar Field cosmology

$$V(\phi) = V_0[(\phi - B)^2 + A] \exp(-\lambda\phi)$$

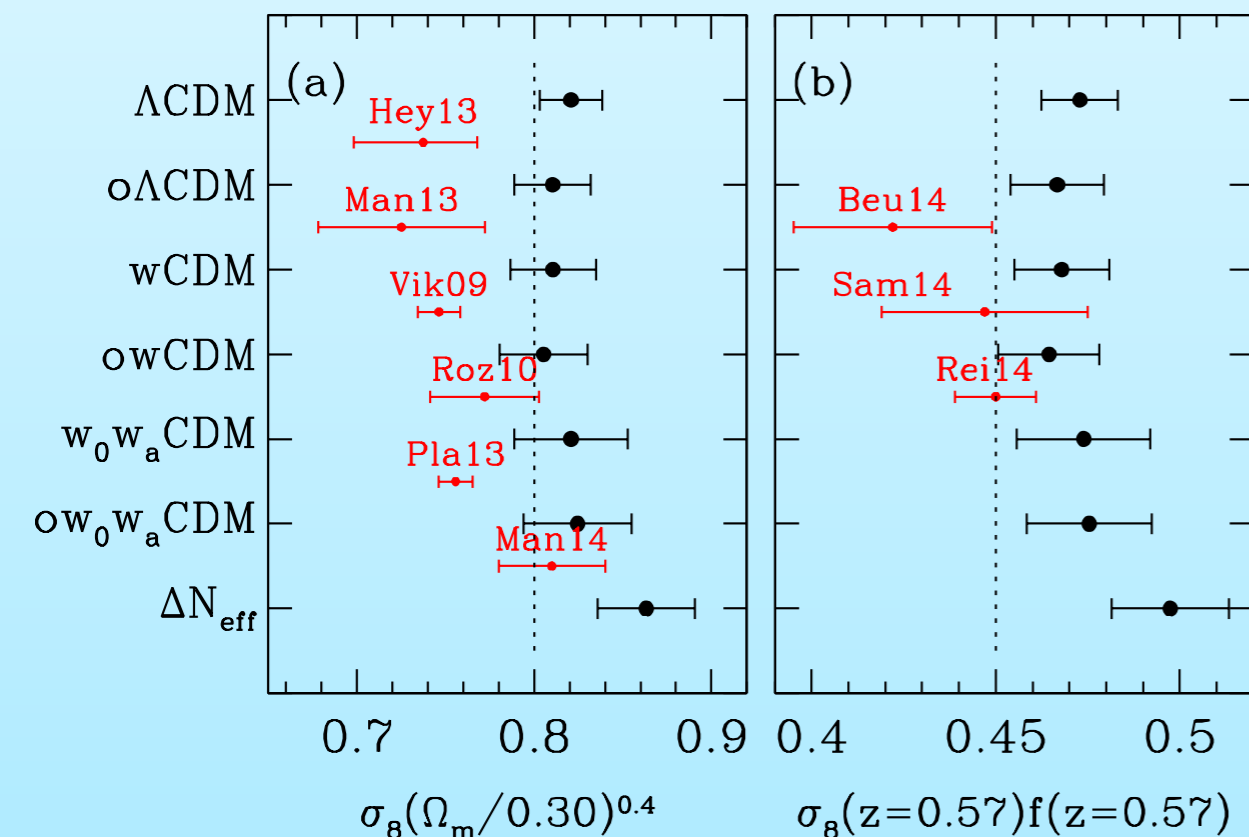
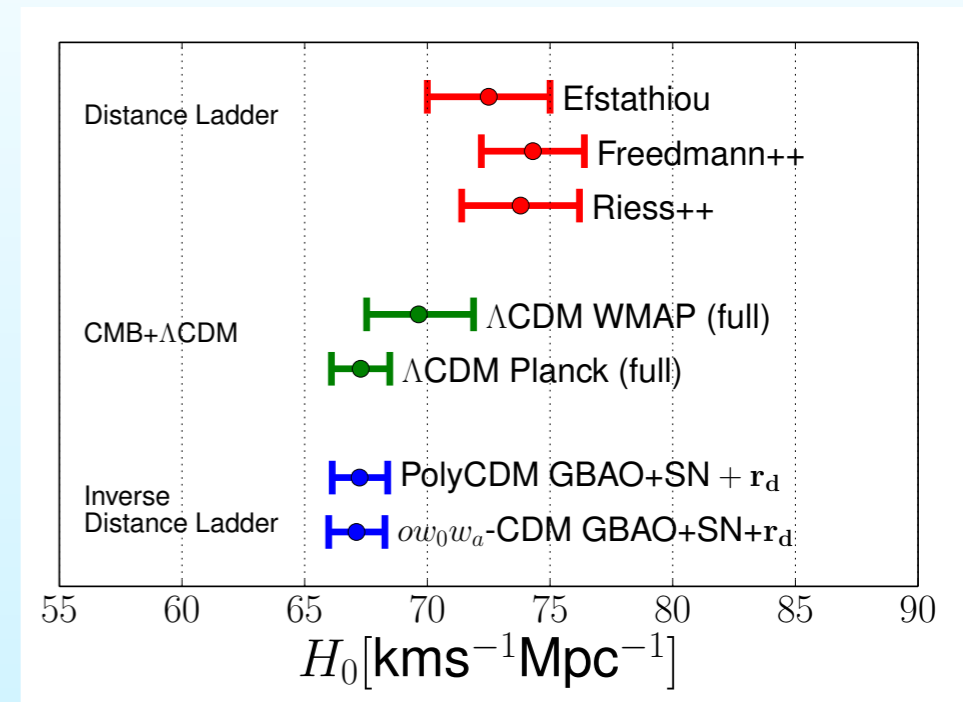


# Scalar Field cosmology



# Scalar Field cosmology

Intriguingly, non-zero  $\Omega_{\text{de}}^e$  with rescaled  $r_d$  reduces tension with distance-ladder measurements of  $H_0$  and with the level of matter clustering inferred from cluster masses, weak lensing, and redshift-space distortions.



Increasing the early dark energy fraction reduces the value of  $\Omega_m$  (see Fig. 13) and will also suppress growth of structure relative to  $\Lambda$ CDM.